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Measurements of suprathreshold hydrogen flux in CASTOR tokamak

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Abstract

A study of the non-equilibrium absorption of hydrogen particles by a plasma-facing absorption probe (AP) made of Group V metals in the CASTOR tokamak environment has been undertaken with the main purposes: (1) to investigate the role of nonmetallic coatings upon plasma-facing materials in D/T inventory and recycling and (2) to develop a method of the registration and diagnostic of the flux of suprathreshold hydrogen. Two series of experiments were performed: with AP of Nb and of V. The absorption of suprathreshold hydrogen particles coming from the tokamak plasma was investigated as a function of AP temperature, AP distance from the plasma, AP bias and of plasma pulse duration. The possibility of a reliable registration of suprathreshold hydrogen was demonstrated in spite of the short plasma pulse duration and a relatively high background H₂ pressure. The composition of the hydrogen flux (molecules, atoms, ions) impinging on the tokamak walls was analyzed, including ion energy distribution.

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1. Introduction

The interactions of neutral and charged hydrogen particles with plasma-facing materials is of prime importance in the present-day fusion devices and their models, as they define D/T fuel recycling and inventory [1,2]. For that reason, measuring of the hydrogen flux from plasma, of hydrogen particle energetic distribution, and of the molecular composition and ion charge, takes on a great interest.

Hydrogen-wall interactions are governed to a great degree by the state of wall surface which conditions the processes of hydrogen absorption/desorption/accumulation. The probability, α , of the absorption of low-energy hydrogen particles (~ 1 to hundreds eV) in metals through a non-metallic film typically covering the metal surface radically depends on the film thickness. In the case of a monolayer

film, this probability is close to that for a clean metallic surface. However, it may decrease dramatically, if the film thickness exceeds one monolayer. In reality, the thickness of nonmetallic coating depends on temperature: at high enough temperatures, typically only one nonmetallic monolayer exists at the surface under vacuum conditions, and hence the probability of absorption of hydrogen particles whose kinetic, internal or chemical energy exceeds 1 eV (suprathreshold hydrogen) is very high. At lower temperatures, both monolayer and polyatomic coatings are possible, and, correspondingly, α may vary over a wide range.

The present work was undertaken to investigate the role played by the nonmetallic coatings on plasma-facing materials in D/T inventory and in the recycling of low energy hydrogen particles in tokamak environments; another goal is developing of a method of suprathreshold particle diagnostics.

A resistively heated absorption probe of Group V metals (niobium and vanadium) is used to trap the flux of hydrogen particles, the amount of which is measured dur-

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56 ing the post-exposure thermal desorption of the absorbed
57 hydrogen.

58 The choice of the probe material is due to a very favor-
59 able combination of such parameters of hydrogen interac-
60 tion with the Group V metals as a large positive heat of
61 hydrogen solution, a high mobility of solute hydrogen
62 depending only weakly on metal temperature, and a high
63 chemical activity of the metal surface giving rise to stable
64 surface films of nonmetallic elements effectively impeding
65 the thermal desorption of solute hydrogen. Due to that,
66 probes of these metals can operate over a wide range of
67 tokamak operation parameters, including high particle flux
68 densities, short pulse duration, and also in devices having a
69 high level of the background molecular hydrogen where
70 other known diagnostic methods may be inapplicable [2,3].

71 The CASTOR tokamak [4] offers a good opportunity for
72 staging of such experiments due to a relatively high, up to
73 10^{20} H/(m² s), flux of suprathreshold hydrogen particles, an
74 easy change of probe samples, and the vacuum system free
75 of carbon-containing impurities. On the other hand, one
76 has to note a relatively high, up to 3×10^{-2} Pa, level of
77 the background hydrogen, and a short, <35 ms, duration
78 of the plasma pulse.

79 2. Experimental

80 Schematic layout of the experimental setup is presented
81 in Fig. 1. A movable resistively heated absorption probe
82 (AP) of the investigated material (0.02 mm niobium foil,
83 or 0.025 vanadium foil) with plasma facing surface area,
84 S_{probe} , ~ 50 cm² is placed in a bakeable vacuum chamber
85 connected to a CASTOR tokamak port through a gate
86 valve. The AP is fixed at movable rod of magnetic drive,

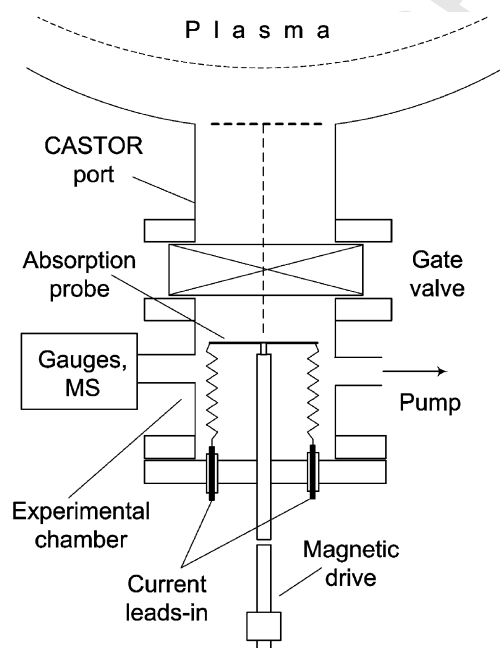


Fig. 1. Experimental layout.

and it can be transported along the system axis over a
 ~ 100 cm distance. AP can be located at different distances
from the tokamak plasma during the absorption experi-
ments, and it can be moved into the experimental chamber
which can be separated from the torus during the Castor
bake-up, as well as for AP conditioning, for the analytical
degassing experiments, experiments on AP surface modifi-
cation, etc.

The AP chamber is pumped by a 200 l/s TMP (40 l/s for
hydrogen). The pressure is measured by a mass-spectrome-
ter and a set of vacuum gauges. The state of AP surface
(coating type and thickness) can be varied *in situ* by con-
trolled reactions with specially introduced chemically
active gases. One can obtain clean Nb(V) surface by heat-
ing the probe to a high temperature. At the first step of this
study, the probes were only heated to moderate tempera-
tures (1900 K for Nb, and 1600 K for V) in vacuum, which
leads to the dissolution of polyatomic oxide films, with
only O monolayer remaining on the surface [5].

Measurements of the absorption of hydrogen from
plasma were performed in two steps. First, the probe was
moved into the CASTOR torus and put at a desirable
distance from the plasma. Then, an operational hydrogen
pressure was established, and a plasma shot was executed.
After the exposure to plasma, AP was moved back into the
separate chamber. The amount of the absorbed hydrogen
was measured at H₂ thermal desorption from the probe
by its heating to 1300 K with the gate valve closed. A
typical example of thermodesorption curve is presented
in Fig. 2, curve 1. The area under this curve is propor-
tional to the number of hydrogen particles absorbed
both from the plasma and from the background molecular
hydrogen.

Second, the above procedure was fully repeated, but
without plasma shot. A corresponding thermodesorption
curve is presented in Fig. 2, curve 2; it refers to the desorp-
tion of hydrogen absorbed only from the background
hydrogen. The number of suprathreshold hydrogen particles
absorbed from the plasma is proportional to the difference
of the areas under the curves 1 and 2 (the hatched area in
Fig. 2).

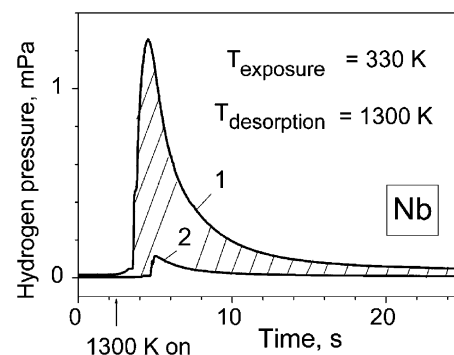


Fig. 2. Hydrogen thermodesorption from the probe at the permanent pumping with a speed 40 l/s for hydrogen: (1) suprathreshold particles and (2) molecular hydrogen.

128 3. Experimental results and discussion

129 3.1. Dependence of the absorption of hydrogen from 130 plasma on probe location

131 The dependence of the amount of hydrogen accumu-
132 lated from the suprathreshold hydrogen flux on probe loca-
133 tion is presented in Fig. 3; it demonstrates a strong
134 dependence of hydrogen accumulation on the distance
135 between AP and plasma. AP position just in line with the
136 tokamak wall is indicated in Fig. 3 as “0” position, which
137 corresponds to a 94 mm distance from the tokamak vessel
138 centre. The “minus” sign refers to an AP shift toward the
139 plasma.

140 The flux of suprathreshold particles consists of both neu-
141 tral atoms and ions. As the concentration of neutral hydro-
142 gen atoms is presumably equidistributed over the whole
143 range of possible AP location, the observed dependence
144 is most probably due to the ions whose flux is different at
145 different distances from plasma because of shielding by
146 the CASTOR limiter, walls, etc., and due to the presence
147 of magnetic field. One can assume that the independence

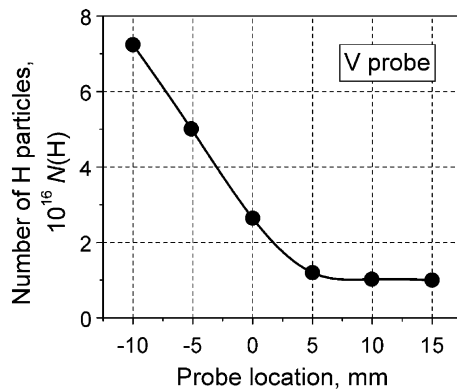


Fig. 3. Dependence of hydrogen accumulation on probe location. $T(\text{exposure}) = 373 \text{ K}$, $T(\text{desorption}) = 1123 \text{ K}$, floating potential.

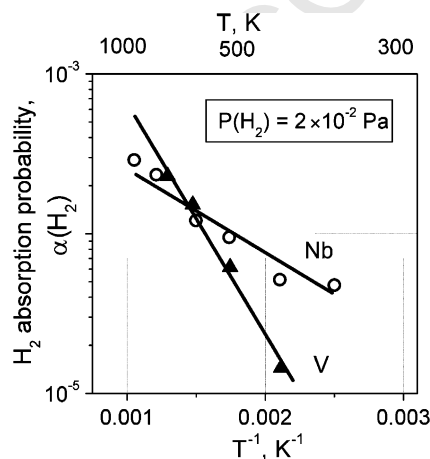


Fig. 4. Temperature dependence of hydrogen molecule absorption probability.

of hydrogen accumulation on probe location (the right
branch of the curve in Fig. 3) indicates at the absence of
an input from ions. Based on this result, the experiments
with neutral particles were performed over this range of
probe location.

3.2. Absorption of hydrogen molecules

One can characterize the state of AP surface with regard
to hydrogen absorption by the dependence of the kinetics
of molecular hydrogen absorption on probe temperature
(Fig. 4). We defined the sticking coefficient of hydrogen
molecules, α_{H_2} , from our experimental data as follows
 $\alpha_{\text{H}_2} = N(\text{abs})/N(\text{inc})$. Here $N(\text{inc})$, the amount of hydro-
gen molecules falling onto the probe surface, equals $K_1 \times$
 $P_1(\text{H}_2) \times t$ where $P_1(\text{H}_2)$ is hydrogen pressure measured
during AP exposure, t is exposure time and K_1 is a propor-
tionality coefficient. The amount of hydrogen molecules
absorbed by the probe, $N(\text{abs})$, equals to the amount of
hydrogen desorbed from AP at thermodesorption, and it
can be experimentally measured as described in Section 2.

As it follows from the data presented in Fig. 4, $\alpha_{\text{H}_2} =$
 $2.1 \times 10^{-2} \exp(-28.2 \text{ kJ/kT})$ for the vanadium probe,
and $\alpha_{\text{H}_2} = 1.5 \times 10^{-3} \exp(-13.6 \text{ kJ/kT})$ for the niobium
probe. These values are typical of an H_2 molecule absorp-
tion by Nb and by V through a surface covered by nonmet-
allic (e.g. O) impurity monolayer resulting from the
specimen annealing under vacuum conditions [5] (note that
 $\alpha_{\text{H}_2} = 0.1\text{--}0.3$ for a clean Nb surface [6]). According to our
experience, that means that the probe surfaces are in a state
giving rise to superpermeation, at least over the investi-
gated temperature range, and hence one can reckon on
an effective absorption of suprathreshold particles.

3.3. Absorption of hydrogen atoms

The number of suprathreshold particles absorbed during
the tokamak discharge, N_{H} , was measured as function of
AP temperature (Fig. 5, curve 3) and of plasma pulse
length (Fig. 6). One can see that the amount of absorbed

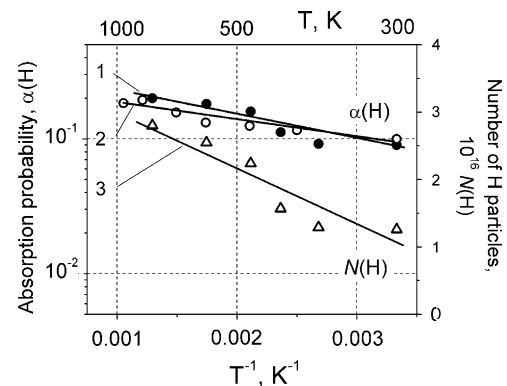


Fig. 5. Temperature dependencies of hydrogen atom absorption probability, (1) vanadium, (2) niobium and of hydrogen accumulation during plasma discharge (3).

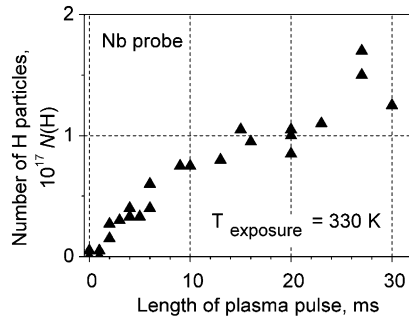


Fig. 6. Dependence of hydrogen absorption on pulse length.

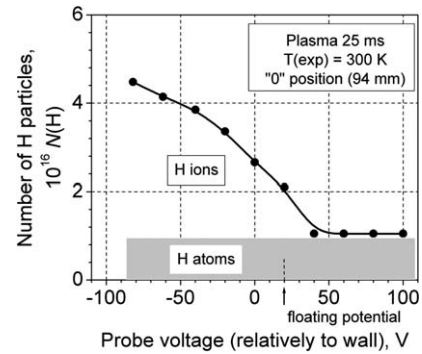


Fig. 8. Dependence of hydrogen accumulation on ion energy.

184 suprathreshold particles depends weakly on probe tempera-
 185 ture and is approximately proportional to pulse length.
 186 The number of hydrogen atoms absorbed during plasma
 187 discharge amounts up to $\sim 10^{17}$ atom H, depending on
 188 probe temperature and on plasma pulse length (Figs. 5
 189 and 6).

190 Concerning the probability of hydrogen atom absorp-
 191 tion, we take its magnitude from the known data on the
 192 implantation coefficient, α_H , at sufficiently high tempera-
 193 tures: when it is guaranteed that the oxygen coating does
 194 not exceed one monolayer. For instance, one can take
 195 $\alpha_H \approx 0.2$ as a typical value for the energy range around
 196 10 eV ($\alpha_H \approx 0.2$ is ascribed to our experimental data for
 197 suprathreshold hydrogen: Fig. 5 curves 1, 2, left axis). As
 198 one could have expected it, the probabilities of hydrogen
 199 atom absorption for niobium (curve 2) and vanadium
 200 (curve 1) probes are very close to one the other.

201 3.4. Bias experiment

202 The study of the absorption of hydrogen ions was car-
 203 ried out with the same procedure as above (Section 2),
 204 but vanadium probe was electrically biased with respect
 205 to the CASTOR walls and was placed at a position where
 206 there is a measurable ion flux (see above). A typical depen-
 207 dence of the ion current on probe bias voltage is presented
 208 in Fig. 7 for two AP positions (see also Fig. 3).

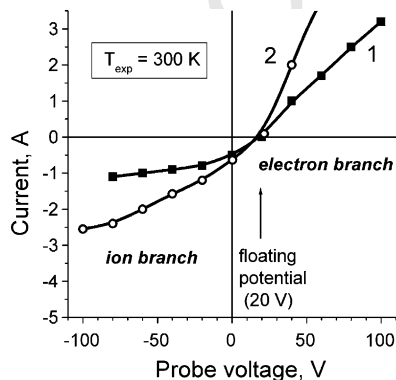


Fig. 7. Dependence of ion current onto the probe on probe bias voltage for two radial positions of the probe: (1) $r_p = 94$ mm ("0" position), (2) $r_p = 84$ mm.

209 The results of the corresponding thermodesorption
 210 experiments (Fig. 8) demonstrate that hydrogen absorption
 211 does not depend on probe bias over the range $+(40\text{--}100)$ V.
 212 One can suggest that such an independence means that the
 213 input of ions into hydrogen absorption is negligible over
 214 this energy range, hydrogen accumulation in the probe
 215 being due only to the input of neutrals.

216 The absorption of neutral hydrogen atoms does not
 217 depend on probe bias and thus it was possible to measure
 218 it separately. The results of these measurements are pre-
 219 sented in Fig. 8 as a gray bar. The difference between the
 220 total desorption from the probe and the desorption indicat-
 221 ing the input of neutral atoms yields the input of hydrogen
 222 ions. It follows from these data that the probability of
 223 hydrogen ion absorption weakly depends on ion energy
 224 over the whole investigated energy range and equals
 225 $\alpha_H = 0.75 \pm 0.05$. These numeric estimates are made on
 226 the presumption that the ion flux consists of H^+ ions [4].

227 Besides, AP being actually a Langmuir probe, its biasing
 228 permits to measure energy distribution of ions in the vicinity
 229 of CASTOR wall.

230 4. Conclusion

231 From a practical point of view, the possibility of a reli-
 232 able registration of suprathreshold hydrogen particles (atoms
 233 and ions) from tokamak plasma with the absorption probes
 234 made of superpermeable niobium and vanadium was demon-
 235 strated in the CASTOR tokamak conditions with its short
 236 plasma pulse and a relatively high background H_2
 237 pressure. The absorption of suprathreshold hydrogen parti-
 238 cles from plasma was investigated as function of probe dis-
 239 tance from the plasma, probe bias, and plasma discharge
 240 duration. These results provide the opportunity to study
 241 composition of the hydrogen flux (molecules, atoms, ions)
 242 impinging upon the CASTOR liner and to evaluate the
 243 fluxes of each component, as well as the energy distribution
 244 of hydrogen ions.

245 Absorption of 7 eV Franck–Condon atoms [4] and of
 246 low energy (up to hundreds of eV) ions by the Group V met-
 247 als, niobium and vanadium, with the surface covered with
 248 oxygen monolayer film (Section 2) was studied from the
 249 viewpoint of physics. It was demonstrated that such a film

250 does not impede the absorption of suprathreshold hydrogen
251 particles, but it suppresses their recycling, which results in
252 an inventory of up to $\sim 10^{17}$ atom H from the plasma, sub-
253 stantially exceeding that from hydrogen molecules.

254 The absorption of suprathreshold hydrogen particles by
255 niobium and vanadium probes was found to be very close
256 to each other. The absorption probability of hydrogen
257 atoms was found to weakly depend on metal temperature
258 over temperature range 300–1000 K. The absorption prob-
259 ability of hydrogen ions does not depend on ion energy
260 over the range 20–100 eV.

261 The measured number of H atoms absorbed during a
262 standard 25 ms discharge is in a reasonable agreement with
263 that estimated from the plasma parameters.

264 The continuation of the present work will be aimed at
265 modifications of the probe surface with the purpose to
266 strongly suppress the absorption of hydrogen molecules
267 and to increase in this way the sensitivity of suprathreshold
268 particle registration. Another promising idea is developing
269 of a method to measure the amount of absorbed hydrogen
270 based on the dependence of probe material resistivity on
271 the concentration of solute hydrogen.

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