INSTITUTE OF PLASMA PHYSICS CZEC: ACAD. SCI.

MODELLING OF THE ENERGY BALANCE IN THE MT-1 TOKAMAK

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Report 21/82

October 1982

Pod vodárenskou věží 4, 182 11 Praha 8

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Plasma processes in the MT-1 tokamak /a=9 cm, R=40 cm, B_T =1.2 T/ were calculated by two numerical codes. The first of these TOKATA /2/ solves a simplified system of equations for the energy balance of electron and ion current components of the plasma and for plasma density. Only heat conductivity is taken into account in the energy balance equations as a main channel for the power losses of the both components. The thermal conductivity coefficient for electrons is approximated by the formula /Alcator scaling/

$$\chi_{p} = 50/n_{p}$$

For the ions, the neoclassical expressions are used

$$\chi_i^{PS} = 170 \frac{Z_{eff}^n_i}{B^2 \sqrt{T_i}} (1 - 1.6 q^2)$$
 collisional region .../2/

$$\chi_{i}^{GS} = 38.6 \frac{q_{i}.T_{i}^{3/2}}{R.B^{2}}$$
 "plateau" region .../3/

$$\chi_i^B = 115.6 \frac{Z_{eff}^n_i}{B^2 \sqrt{T_i}} \cdot q^2 \cdot \left(\frac{R}{r}\right)^{3/2}$$
 banana regime.../4/

/ units $\chi_{i}[cm^{2}ms^{-1}]$, $n_{i}[10^{13}cm^{-3}]$, $T_{i}[eV]$, B[kG]/

and the following common formula is included the TOKATA code

$$\chi_i = AHI * \left(\chi_i^{PS} + \frac{1}{\sqrt{\chi_i^B} + \sqrt{\chi_i^{GS}}}\right) \dots /5/$$

where constant AHI allows to increase artificialy the power losses from ion component by thermal conductivity.

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The resistivity of the plasma is supposed to be classical /Spitzer/, with possibility of its enhancement by a constant AR

$$\eta = AR \times \eta_{class}$$
... /6/

Results of the preliminary calculation for a regime with the current I_p = 20 kA and value of the density at the centre of plasma $n_e(0)$ = 3. 10^{13}cm^{-3} /see fig. 1/ have shown, that the coputed central values of T_e and T_i are in reasonable agreement with experimental data for AHI = 3, AR = 1 and Z_{eff} = 3. But the calculated radial profiles are probably too flat, because of neglecting radiation, diffusion and charge-exchange energy losses which play a dominant role at the pherifery of the plasma column.

The second code TOKSAS calculates the detailed energy balance of the ion component so that the energy losses by diffusion $P_{\mbox{dif}}$ and by charge-exchange processes $P_{\mbox{ex}}$ are taken into account, too.

$$\frac{dQ_i}{dt} = P_{ei} - P_{cond} - P_{dif} - P_{ex} \qquad ... /7/$$

The neoclassical values of thermal ion conductivity are the same as in TOKATA code, but common formula differs a little

$$\chi_{i} = \frac{\chi_{i}^{\mathcal{B}} \left(\chi_{i}^{\mathcal{P}S} + \chi_{i}^{\mathcal{G}S}\right)}{\chi_{i}^{\mathcal{P}S} + \chi_{i}^{\mathcal{G}S} + \chi_{i}^{\mathcal{B}}} \qquad ... /8/$$

The radial profile of the neutral atoms /see fig.2/ used for enumerating of the diffusion and charge exchange losses P_{ex} is calculated independently by solving Boltzman transport equation [4], supposing fixed profiles $T_i(r)$, $T_e(r)$, $N_e(r)$

and density of the neutral atoms at the scrape-of-layer $n_{\rm g}(a)$ as input parameters.

From the main numerical result - the radial profile of ion temperature $T_i(r)$ /see fig. 3/, the radial dependence of all power channels can be enumerated as it is shown on the fig. 4 a, b.

Another two important quantities, suitable for direct comparison of the computed values with the experimental data, are calculated by the code TOKSAS.

One of them the energy spectrum of the fast charge-exchange atoms /see fig. 5/, escaping from the plasma column along the central chord is evaluted by expression

$$\left(\frac{dJ_a}{dE}\right) = \frac{1}{4\pi} \int_{-a}^{a} n_a(r) \cdot n_i(r) \cdot \gamma_i(r, E_0, T) \cdot \sigma_{Ex} \cdot \mathcal{N}(E_0) \cdot e^{-\tau} dr \dots (g)$$

where ψ_i - distribution of plasma ions /supposed to be maxwellian /

Gex-charge exchange cross-section

 τ - factor, describing attenuation of the fast neutral flow by the plasma

Moreover the radial distribution of photons emitted by excited hydrogen atoms can be calculated by subroutine HALPH /see fig. 6/

$$J_{pq} = \frac{1}{4\pi} \cdot n(p) \cdot A_{pq} = \frac{1}{4\pi} p^2 r(p) \cdot n_a \cdot A_{pq} \qquad \left[\frac{photons}{cm^3 s. stree} \right] \dots (10)$$

where p, q - principal quantum numbers for the given spectral line

Apq - transition probability

 $r_1(p,n_e,T_e,n_a)$ - coefficient of the proportionality between excited and ground state atoms in the collisional - radiative model [5] approximated for our purporses as

$$r_1(3) = 1,13.10^{-4} \ln(0.78 \cdot 10^{-13} + n_e + 1)$$
 $H_{\alpha} \dots /11/$
 $r_1(4) = 1,8.10^{-5} \sqrt{n_e \cdot 10^{-13}}$ $H_{\beta} \dots /12/$

For direct comparison with optical diagnostic data, the line-integrated intensity of H_{α} and H_{β} is calculated

$$I_{pq}(h) = \int J(x) dx \qquad x = \sqrt{r^2 - h^2} \qquad/13/$$

$$-\sqrt{a^2 - h^2} \qquad I_{pq} \left[\frac{photons}{cm^2 \cdot s \cdot strd}\right]$$

References:

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- [3] Bernard M.: REPORT EUR-CEA-FC-891, Avril 1977, FONTANAY-AUX-ROSES
- [4] Izvoschikov A.B., Petrov M.P.: Fizika plazmy. 2(1976)2, p.212

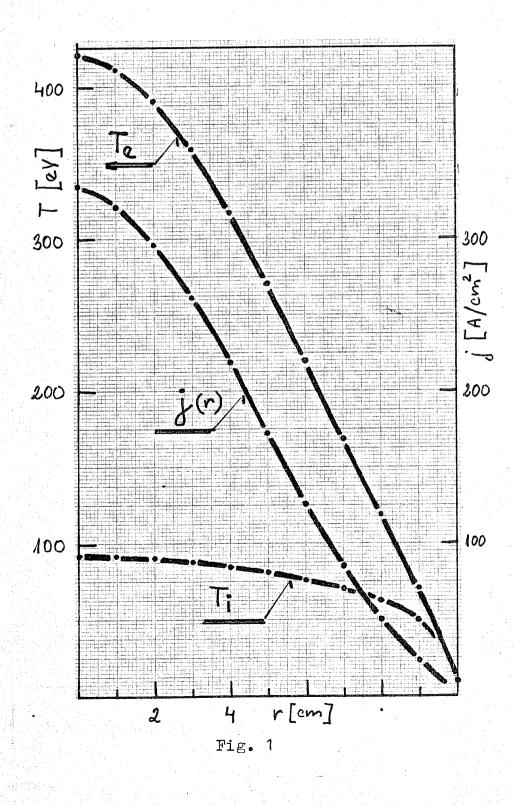
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 4(1978)6, p.1205
- [5] Jonson L.C., Hinnov E.: Princeton University Report MATT 909,
 August 1972

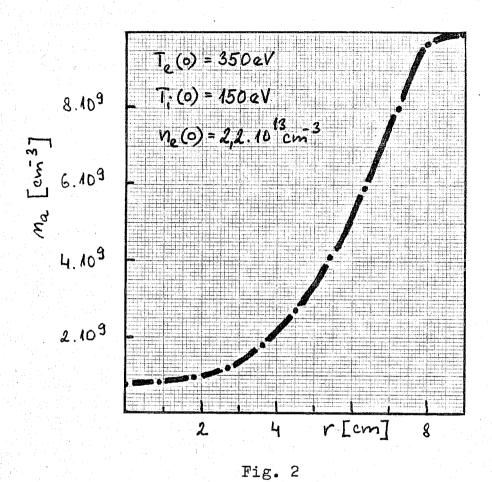
TOKATA input parameters:

$$a = 9$$
 cm, $R = 40$ cm, $B_T = 1.2$ T, $I_p = 30$ kA, $q(a) = 4.05$

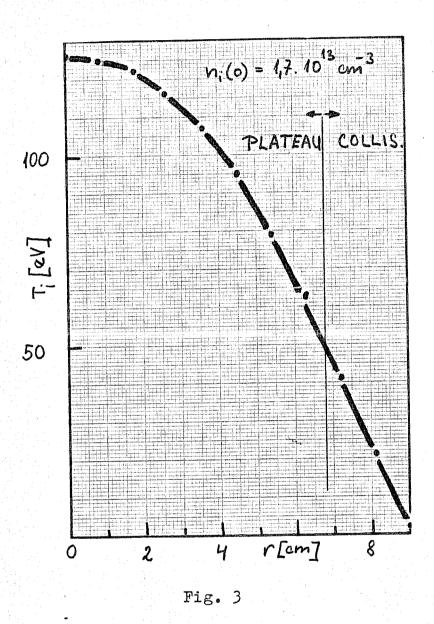
TOKATA output parameters:

$$U_z = 2.2 \text{ V}, \text{ q(0)} = 1.43, Q_p = 74 \text{ J}, P_{OH} = 66 \text{ kW}$$
 $T_E^* = 1.12 \text{ ms}, \beta_p = 0.44, p = 770 \text{ J/m}^3$

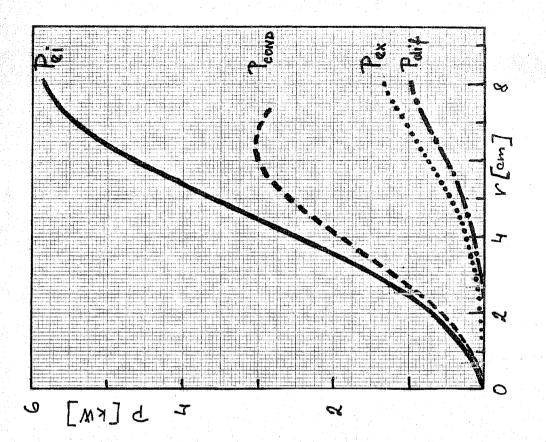




Radial density profile of hydrogen atoms calculated by the code TOKSAS. The profiles of density and temperature of electrons and ions, the density and temperature at neutrals at the boundary were fixed.



Radial profile of ion temperature.
Using the density profile of neutrals calculated above and assuming the electron temperature profile to be known, the code TOKSAS calculates the ion temperature profile.



Tig. 4b
Integrated energy gain and losses
of ions $R(r) = 2\pi R$. $\int 2\pi r' \, Q(r') \, dr'$

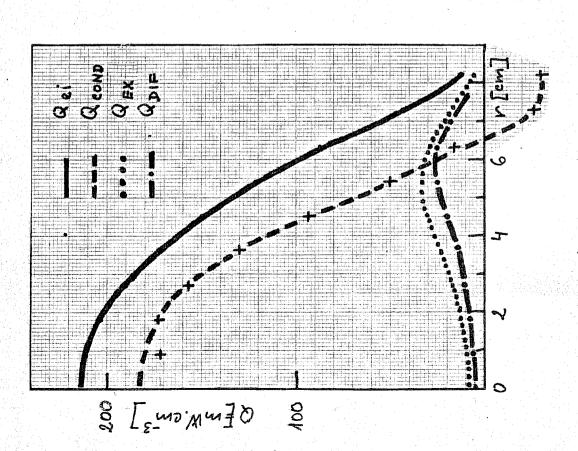


Fig. 4a Energy gain and different types of energy losses of ions a unit volume

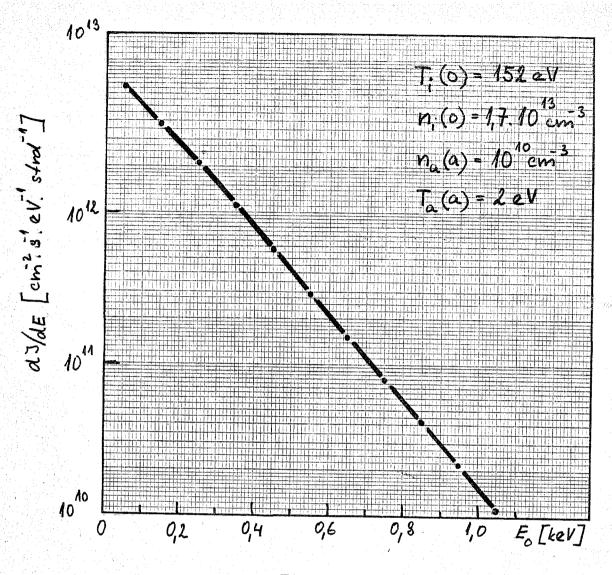


Fig. 5

Energy spectrum of the fast hydrogen atom flow

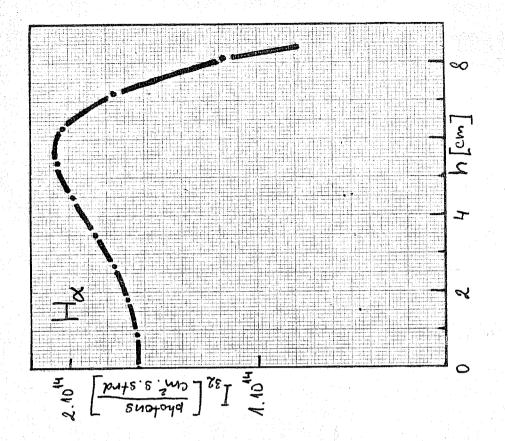


Fig. 6b Intensity of H_{α} , line-integrated along different chords.

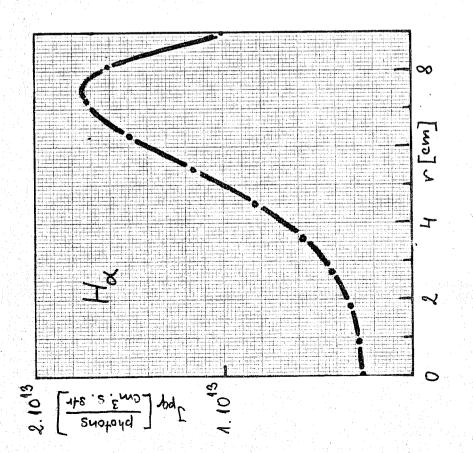


Fig. 6a Volume intensity of $H_{\mbox{\scriptsize c}}$